

Efficiency of Photochemical Stages of Photosynthesis in Purple Bacteria (A Critical Survey)

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Abstract—Based on currently available data, the energy transfer efficiency in the successive photophysical and photochemical stages has been analyzed for purple bacteria. This analysis covers the stages starting from migration of the light-induced electronic excitations from the bulk antenna pigments to the reaction centers up to irreversible stage of the electron transport along the transmembrane chain of cofactors-carriers. Some natural factors are revealed that significantly increase the rates of efficient processes in these stages. The influence on their efficiency by the “bottleneck” in the energy migration chain is established. The overall quantum yield of photosynthesis in these stages is determined.

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The world of photosynthesis is a unique symbiosis of processes responsible for the annual reservation on Earth of about $3 \cdot 10^{24}$ J of energy in the biomass of plants, algae, and photobacteria [1, 2]. Historically, Krasnovsky's reaction was the first model of photosynthesis: Krasnovsky realized in the liquid phase the reversible photoreduction of chlorophyll *a* molecules with ascorbic acid that is accompanied by an increase in the free chemical energy of the system [3-5]. It should be noted that Krasnovsky was inspired to study photochemical reactions of chlorophyll by Academician A. N. Terenin, who in 1943 discovered the relatively long-lived triplet state in molecules of dyes and developed ideas about their effective involvement in photochemical reactions [6]. In fact, the first works of Krasnovsky and his colleagues revealed that it was the triplet state of chlorophyll that was responsible for its photoreduction in solutions (works [3-5] and references therein). Later, reactions of the chlorophyll triplet state were studied in detail using a flash-photolysis approach ([7] and references). The works of Krasnovsky and Terenin [3-6] had great resonance in science, and many photochemists replenished the series of researchers in photosynthesis. However, over time some works

appeared whose results did not fit the triplet mechanism of photosynthesis. The study of Dmitrievsky et al. [8] stimulated by A. N. Terenin showed that the lifetime of chlorophyll *a* in algae and chloroplasts is approximately fourfold lower than the chlorophyll *a* lifetime in solutions. Similar results were obtained by Brody and Rabinovich for the algae *Chlorella* [9]. However, for a long time a burning question existed for researchers: why are no triplets found in chlorophylls in preparations of photosynthesizing organisms? After numerous unsuccessful attempts in laboratories of different countries, triplets were recorded in chlorophyll of chloroplasts, but the yield was extremely low [10]. This effect was, in particular, mainly determined by an exclusively high rate of energy utilization within the reaction centers of photosynthesis, as first established in the works of Godik and the author of this article [11-13]. These works have shown that in four purple bacteria the fluorescence lifetime of bacteriochlorophyll (BChl) is in the range of dozens of picoseconds. This lifetime increased by nearly an order of magnitude during photooxidation (outcome from the active state) of the reaction centers (RC). Later the picosecond range of the singlet BChl lifetimes *in vivo* was confirmed in studies of Paschenko et al. [14], Campillo et al. [15], and in some subsequent works using picosecond laser approaches. Based on these ultrashort times, the principle of instantaneous action of photosynthesis was formulated

Abbreviations: BChl, bacteriochlorophyll; RC, reaction center; SEE, singlet electronic excitation; * indicates the SEE presence in a molecule or in a group of molecules.

22] seems to be fitting. According to this theory, statistics of SEE transfers from an excited donor molecule (D^*) onto an acceptor molecule (A) in general is described by the following formulas:

$$\mathbf{k}_{D,A} = (\mathbf{t}_{D,A})^{-1} \sim f(\varphi_{AB}, \psi_A, \psi_D)^2$$

$$(\mathbf{R}_{cr}/L_{D,A})^6 n^{-4} \tau_R^{-1} \int F_D(\nu) \varepsilon_A(\nu) \nu^{-4} d\nu, \quad (1)$$

$$(\mathbf{R}_{cr})^6 \sim 2/3 (n^{-4} \tau_R)^{-1} \int F_D(\nu) \varepsilon_A(\nu) \nu^{-4} d\nu, \quad (2)$$

where $\mathbf{k}_{D,A}$ is an averaged rate constant of the SEE transfer from D^* to A; $\mathbf{t}_{D,A}$ by determination is the average time of this transfer; $f(\varphi_{AB}, \psi_A, \psi_D)^2 = [\cos\varphi_{AB} - 3\cos\psi_A \cos\psi_D]^2$ is a factor of mutual position of dipole transition moments of the D^* and A^* molecules on their transitions between the basal and lower excited states; φ_{AB} , ψ_A , and ψ_D are angles between these dipoles [21]; \mathbf{R}_{cr} is the critical distance of the SEE migration from D^* onto A at the constant $\mathbf{k}_{D,A}$ value equal to the sum of intramolecular constants of SEE deactivation in D^* molecules, the \mathbf{R}_{cr} value being determined in chaotic three-dimensional ensembles of molecules with mean value of the factor $\xi(\varphi, \psi, \theta)^2 = 2/3$ [21]; τ_R is the radiation lifetime of BChl; $L_{D,A}$ is the distance between centers of chromophores of the donor and acceptor molecules; n is the light refraction coefficient in the medium; τ_D is the SEE lifetime (or fluorescence) of the donor molecules; $F_D(\nu)$ and $\varepsilon_A(\nu)$ are the normalized fluorescence spectrum of the donor and the absorption spectrum of the acceptor; according to the Einstein's law [27], the expressions $\int F_D(\nu) d\nu/\tau_r \nu^3$ and $\int \varepsilon_A(\nu) d\nu/\nu$ are proportional to oscillator forces, or otherwise, to dipole transition moments for singlet transitions $D^* \rightarrow D$ and $A \rightarrow A^*$; ν is an optical frequency in cm^{-1} . Note that the formulas (1) and (2) were deduced on the assumption that the distances between centers of the dipoles $D \rightarrow D^*$ and $A \rightarrow A^*$ were significantly greater than their effective lengths.

We will analyze below the specifics of all factors included in the basic formula (1) *in vivo*. We will compare the SEE migration rates in purple bacteria with the migration in chaotically located BChl monomers at similar distances between the chromophores.

Orientation factor $f(\varphi_{AB}, \psi_A, \psi_D)^2$. The dipole transition moments (further in the text referred to as "dipoles") of BChl molecules involved in the circular LH1 complexes are located within membranes nearly in a single plane [28, 29]. The mean value of this factor for two-dimensional chaotic ensembles of molecules is known to be higher than for three-dimensional chaotic ensembles (2/3). Therefore, in purple bacteria in the ensembles of B875 antenna dipoles the fraction of disadvantageous positions decreases with respect to dipoles of the P870 special pairs of RC. Because of ultrafast energy migration along the B875 ring, the SEE from molecules with the

dipoles situated unsuccessfully relative to P870 enter virtually instantaneously into B875 molecules with the successfully located dipoles, which optimizes the migration from the B875* complex onto the P870 pair.

Factor $L_{D,A}$. As mentioned above, distances between the centers of chromophores of the B875 antenna molecules and the P870 special pair of RC are very great and reach 45 Å [25, 26]. However, there is also a positive consequence: all molecules B875 in the LH1 structure are able to directly deliver SEE onto the P870 special pair located in the center of their circle.

Factor n^{-4} . The use by Förster of the light refraction coefficient n in formulas (1) and (2) was quite adequate for homogenous liquid and solid solvents of dyes that were used in studies of energy migration. However, this parameter does not have a physical sense in membranes of dozens of Å thickness. Therefore, in the case of photosynthesis, n^{-4} in formulas (1) and (2) has to be replaced by the dielectric permeability coefficient ($\chi^2 = n^4$) [30, 31] in the nearest microenvironment of the donor and acceptor molecules. According to numerous literature data, in different proteins the factor χ is ~ 3.0 [32]. Obviously, for the SEE *in vivo* the migration values χ are important in the case of location of BChl molecules in a particular medium and on the time intervals specific for separate migration acts, i.e. on frequencies where only the small electron mass can react to rapid changes in the fields [33].

The interior of the membrane containing in purple bacteria the B875 and P870 molecules consists of approximately equal transmembrane parts of proteins and phospholipid tails. The specificity of the membrane protein primary structures discovered by Prof. Zuber is extremely important: their transmembrane α - and β -helices contain virtually no aromatic and polar amino acids (they are concentrated in the terminal chains that extend from the membrane) [34, 35]. The phospholipid tails forming another component of the mass of membranes also are very hydrophobic. Thus, the interior of membranes mainly consists of nonpolar lipidic groups. Correspondingly, their electrons on single chemical bonds have minimal ability for polarization.

In the reference book [36], the light refraction coefficient (n) values for most nonpolar chemical compounds are within $1.28 \leq n \leq 1.31$, and from the formula $\chi = n^2$ it follows that their minimal values of χ are very small, $1.64 \leq \chi < 1.72$.

In our work [31] the χ value was evaluated for hydrophobic fragments of RC polypeptides around the P870 special pair. On consideration of small amounts of bound water, for the minimal χ of the medium the χ factor in the internal layer of about 20 Å with the gradient from the center to periphery was found to be in the limits of 1.65-1.78 [31]. Considering the influence of the P870 molecules themselves, for the more likely parameter of the medium $\chi = 1.72$ this value occurs in the limits of 1.65-2.00.

Optical frequencies. The optical frequency (ν^{-4}) plays an important role in formulas (1) and (2). On comparison with BChl B800 monomers *in vivo*, the frequency of the long-wavelength band in B875 molecules is 9% lower that under other equal conditions increases approximately by 42% the migration rate constant according to formula (1). On comparison with typical dyes such as acridine, fluorescein, rhodamine, etc., this gain is about 400%.

Bacteriochlorophylls seem to be the red-most natural dyes. In a neutral organic medium, the long-wavelength absorption band of BChl-*a* has a maximum near 770 nm. It is important that it can be easily shifted to the "red" region of the spectrum. Even BChl-*a* monomers in the antenna complexes of LH2 have peaks at about 800 nm due to bonds with the carrying polypeptide chains. The BChl-*a* pair complex in LH2 has peaks at about 850 nm and in LH1 at about 875 nm, also due to the bond with the carrying polypeptides plus strong bonds in the pairs and with neighboring pairs in the circular structures. Such a train of absorption bands enlarges the spectrum of active absorption of solar energy and brings its red boundary in purple bacteria up to 910-920 nm on the basis of BChl-*a* and up to 1030-1040 nm on the basis of BChl-*b*. This boundary is rather close to the optimal boundary for all energy systems of the photoelectric type. Upon the selective absorption of solar radiation by the atmosphere, the optimal boundary on the surface of our planet is close to 1050 nm [37, 38].

Spectral overlap integral $\int F_D(\nu) \epsilon_A(\nu) \nu^{-4} d\nu$. Chromophores of the P870 pairs are located closer to each other than chromophores of α/β pairs of B875, which results in a virtual coincidence of the 0-0 transitions in their long-wavelength absorption bands. Consequently, the overlap degree of the B875 fluorescence and the P870 absorption spectra in formula (1) reaches 87% of the maximum, resulting theoretically in matching the peaks of these spectra [39].

Excitons. At the above-indicated distances between the centers of neighboring chromophores of molecules of P870 (7.5 Å) and B875 (9 Å) [26, 31], values of their interaction energies reach ten meV. In such associates excitons arise: in the RC the SEE includes a pair of P870 and in the LH1 chains it includes 4-6 neighboring molecules of B875 [40, 41]. It is known from physics that in associates including z identical interacting molecules, the long-wavelength absorption band is split into z separate bands. And the sum of dipoles of these bands is retained, i.e. it is approximately equal to the sum of dipoles of z monomers of the corresponding molecules. In the LH1 complexes of purple bacteria, because of the circular symmetry of B875 molecules the majority of these bands are forbidden, and thus only two bands with the wave vectors $k = \pm 1$ remain [40, 41]. The mutual position of dipoles in the neighboring molecules of P870 and B875 provides for the domination of the long-wavelength

bands. According to calculations, at real mutual spectral positions of the absorption bands in the B875 complexes and P870 pairs the ratio of the forward and reverse constants of SEE migration in the region of B875* \leftrightarrow P870 has to increase $3 < \delta < 4$ times [42]. In special pairs of RC from the bacterium *Rhodobacter sphaeroides*, the long-wavelength band near 870 nm is occupied by approximately 85% of the strength of the sum of dipoles of two molecules. On SEE migration, the BChl associates with delocalized exciton can be considered as individual molecules. This idea is confirmed by experimental data of Razjivin et al. [43] that in the *Rhodospirillum rubrum* chromophores, after the exciting pulse the B875 band bleaching calculated per absorbed photon is 4-5 times greater than the bleaching corresponding to the BChl monomer. The exciton effect also significantly increases the SEE migration rates between B875* and P870, because, according to the theory of this phenomenon [21, 22], their amplitudes are proportional to the product of dipoles of the interacting molecules.

Comparison by formula (1) of the migration rates between chaotically distributed in space BChl B800 monomers in protein medium and between the B875 and P870 complexes *in vivo* indicates that the total gain from all the above-listed factors for values of these constants at SEE migration from B875* onto P870 is 100-200. Because of this, the efficient constant value of the SEE migration from the B875 complexes onto the P870 pairs reaches *in vivo* several dozens of inverse picoseconds [44].

COUPLING OF MIGRATION WITH SUBSEQUENT ENERGY TRAPPING AND CONVERSION IN REACTION CENTERS. GROSS QUANTUM YIELD

The ultrafast migration of SEE from the B875* antenna molecules onto the P870 pairs in the RC still does not solve the problem of efficient photosynthesis, because SEE can be returned from P870* onto B875 through the same mechanism. Let us perform approximate calculations based on experimental data on the SEE migration and trapping in purple bacteria.

Data of Godik [45] and Katiliene et al. [46] have shown that the rate of trivial SEE deactivation in the B875 antenna complexes is $\sim(650 \text{ ps})^{-1}$. To provide the quantum yield of SEE trapping in the RC to be about 90%, the rate of this trapping has to be at least an order of magnitude higher, i.e. $\sim(65 \text{ ps})^{-1}$. Considering that 16 pairs of B875 correspond to each P870 pair in the RC, let us assume that the rate of reverse migration of the SEE from P870* onto B875 can be 16-fold greater, i.e. $\sim(4 \text{ ps})^{-1}$. Up to now, in many reviews it is stated that the rate of the SEE primary trapping in the RC of purple bacteria is $\sim(3 \text{ ps})^{-1}$. Such rate of SEE trapping and the ratio 1 : 16 of the rates of forward and reverse migrations result

in the quantum yield (QY) of the trapping to be ~ 0.65 , even on assuming the irreversibility of this trapping.

This discrepancy is significantly removed due to above-mentioned features of excitons [40], which increase 3–4 times the ratio of migration rates, and to new kinetic data for RC obtained in the pioneering works of Shuvalov's laboratory.

In the works of Shuvalov [47, 48], the RC state was recorded with the transfer of an excited electron onto P800 with a small loss of energy and a very short lifetime of ~ 1 ps. In subsequent works of Shuvalov et al. [49, 50], the first excited state of the RC was detected within ~ 0.2 ps. It seemed to be charge transfer complexes (CTC) in the P870 pairs in which noticeable electric dipoles were produced. In the field of such a dipole, a specific water molecule was polarized between the P870 dipoles [51], along with quenching up to 0.02 eV of the SEE energy [52, 53].

Based on the above-presented kinetic data on the SEE migration in the LH1 complexes and their trapping in the RC of purple bacteria, a model was designed describing the succession of the migration primary processes and energy trapping in the LH1 complexes of purple bacteria. From this model (presented below), the quantum yield value of photosynthesis was determined for these organisms.

MODEL OF PRIMARY PROCESSES FOR PURPLE BACTERIA

In the model presented in Fig. 2, the electron transfer from H760^- onto Q_A with constant $k_{e3+} \approx (200 \text{ ps})^{-1}$ is virtually irreversible because the energy difference between these states is close to 0.55 eV [52, 53]. The subsequent electron transfers from Q_A^- onto the quinone pool dissolved in the membrane occur within dozens of milliseconds. These transfers are determined by the rate of quinone diffusion and limit the real carrying capacity of the electron transport chain in purple bacteria. The charge deficiency arising in the special P870^+ pair upon the electron transfer is compensated by the RC-bound

molecule of a heme-containing protein, cytochrome *c* ($\text{P870}^+ + \text{Cyt-c} \rightarrow \text{P870} + \text{Cyt-c}^+$), and within approximately two microseconds after that the RC is able to function again.

All relay-race transfers of SEE and electrons occur in BChl complexes fixed in the membranes, and they therefore can be modeled using first order balanced differential equations. The model for the SEE migration processes and their conversion in the purple bacteria RC is presented below.

$$d\{\text{B875}^*\}/dt = \delta k_m \{\text{P870}^*\} - (k_m + k_{wA}) \{\text{B875}^*\} \quad (\text{SEE migration}), \quad (3a)$$

$$d\{\text{P870}^*\}/dt = k_m \{\text{B875}^*\} + k_{\text{CTC-}} \{\text{CTC}\pm\}^* - (k_{\text{CTC+}} + \delta k_m + k_{w\text{P870}}) \{\text{P870}^*\}, \quad (3b)$$

$$d\{\text{CTC}\pm\}^*/dt = k_{\text{CTC+}} \text{P870}^* + k_{e1-} \{\text{P800}^-\} - (k_{\text{CTC-}} + k_{e1+} + k_{w\text{CTC}}) \{\text{CTC}\pm\}^*, \quad (3c)$$

$$d\{\text{P800}^-\}/dt = k_{e1+} \{\text{CTC}\pm\}^* + k_{e2-} \{\text{H760}^-\} - (k_{e2+} + k_{e1-} + k_{w\text{P800}}) \{\text{P800}^-\}, \quad (3d)$$

$$d\{\text{H760}^-\}/dt = k_{e2+} \{\text{P800}^+\} - (k_{e3+} + k_{w\text{H760}}) \text{H760}^- \quad (3e)$$

$$d\{\text{Q}_A^-\}/dt = k_{e3+} \{\text{H760}^-\} \quad (3f)$$

In these primary stages, all processes are 99.99% finished within 1.5 ns; therefore, the influence of slower processes of electron transfer in the subsequent stages $\text{Cyt-c} \rightarrow \text{P870}$ and after $\{\text{Q}_A^-\}$ can be neglected. On the modeling, SEEs were initially given to the B875 and P870 complexes in the ratio of 16 : 1. The major fraction of SEEs from B875* successfully reaching P870 was converted within it into separate charges with different sign. After three stages of the transfer, electrons dislodged from P870 in the state $\{\text{CTC}\pm\}^*$ reached bound ubiquinone

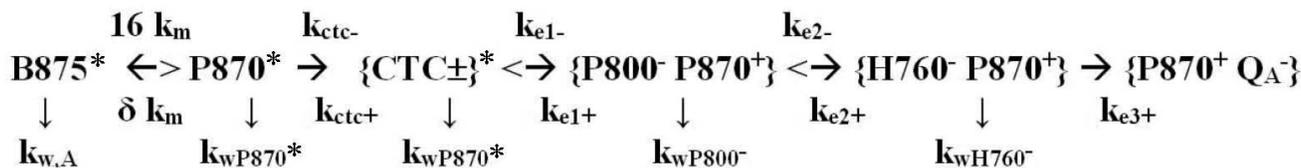


Fig. 2. Scheme of primary processes in purple bacteria. B875^* is the excited state in the B875 antenna molecules that can migrate onto a special P870 pair of an RC with averaged constant $\delta k_m \approx (35 \text{ ps})^{-1}$ [17, 40, 42] at the reverse constant $16 k_m$. In P870^* the state $\{\text{CTC}\pm\}^*$ (abbreviation for Charge Transfer Complex) is produced with the constant $k_{\text{CTC+}} = (0.2 \text{ ps})^{-1}$ [47, 48]; $\{\text{P800}^- \text{P870}^+\}$ is the RC state after the first electron transfer onto P800 with constant $k_{e1+} = (3 \text{ ps})^{-1}$ [45, 46]; $\{\text{H760}^- \text{P870}^+\}$ is the RC state after the secondary transfer of the electron onto bacteriopheophytin with the constant $k_{e2+} = (1 \text{ ps})^{-1}$ [45, 46]; $\text{P870}^+ \text{Q}_A^-$ is the RC state after the third, irreversible electron transfer with constant $k_{e3+} = (200 \text{ ps})^{-1}$ [2, 4, 42]; $k_{\text{CTC-}}$, k_{e1-} , and k_{e2-} are rate constants of the reverse transfers; k_{wA} , $k_{w\text{P870}^*}$, $k_{w\text{CTC}}$, $k_{w\text{P800}}$, and $k_{w\text{H760}^-}$ are rate constants of deactivations of the corresponding states in the RC.

{ Q_A }. Their fraction rather accurately corresponds to the value of effective quantum yield in purple bacteria.

Now it is necessary to introduce the correction factors and values of some constants. The correction by average values of the factors $f(\Phi_{AB}, \Psi_A, \Psi_D)^2$ *in vivo* and in chaotic ensembles of molecules $0.95 : 2/3 = 1.425$. The correction is performed by the absorption peaks of B875 and B800: these peaks in *Rhodobacter sphaeroides* at 810 and 864 [50] are given by the ratio $(864/810)^6 = 1.443$ [54].

The modeling has shown that the most important constant is k_m . The distances between the centers of B875 and P870 chromophores are 45.3 \AA [25, 26]. At such distance it is reasonable to use Förster's fundamental theory of inductive resonance. Let us write formula (1): a) for our case, and b) in a special form for the chaotic ensemble B800, when $L_{D,A} = R_{cr}$ and $\tau_D = \tau_R$. Then for (b) by the definition: $k_{D,A} = (\tau_R)^{-1}$ and $\xi_{D,A}(\psi, \theta_1, \theta_2)^2 = 2/3$. The division each against the other, respectively, the left and right parts of these factors with the correction factors, will result in the known Förster expression:

$$k_{D,A} = (t_{D,A})^{-1} = 1.425 \times 1.443 \frac{\xi_{D,A}(\psi, \theta_1, \theta_2)^2}{2/3 \tau_R} (R_{cr}/L_{D,A})^6 \cdot (4)$$

For the major spectral fraction of the purple bacteria pigment BChl-*a* B875, the R_{cr} values on the SEE migration in photosynthetic membranes were found to be 80 \AA [55] and 78.8 \AA by an independent method [31]. Their mean value for the wavelength $875 \pm 5 \text{ nm}$ is $R_{cr} = 79.5 \pm 2 \text{ \AA}$. From the literature the radiative lifetime of BChl is known to be $\tau_R = 18 \pm 2 \text{ ns}$ [56]. The effective value $\xi_{D,A}(\psi, \theta_1, \theta_2)^2$ for case (a) is close to 0.95, because due to the ultrafast migration along the B875 ring the SEE rapidly reaches the B875 α/β -pairs located optimally relative to P870. Place these data formally into formula (4):

$$k_{D,A} = (t_{D,A})^{-1} = \frac{0.95 \times 1.425 \times 1.443}{2/3 \times 18 \times 10^{-9}} (79.5/45.3)^6 \approx 4.74 \times 10^9 \text{ s}^{-1}. (4a)$$

During the SEE migration, let us use the exciton factor described in the first part of this work according to which the ratio of the forward and reverse constants of B875 \leftrightarrow P870 migration is $16 : \delta \approx 4.6$. Taking into account that the strength of the long-wavelength dipole P870 $\approx 84\%$ of the sum of dipoles of two BChl molecules producing P870 [54], this factor in them is to be increased 1.68-fold. Similarly for excitons involving $4-6 \approx 5$ B875 molecules, this factor is approximately equal to 4 [24, 43]. These factors together increase 7-fold the values of the migration constants δk_m and $6 k_m$ in Eqs. (3a) and (3b), which results in the following evaluation:

$$k_{D,A} = (t_{D,A})^{-1} = 33.2 \times 10^9 \text{ s}^{-1}. (4b)$$

However, under the actual redox potentials of the medium *in vivo*, the fraction of photoactive RCs is 85-90% (the other 10-15% inactive P870 quench the SEE with rate of $(40-50) \times 10^9 \text{ s}^{-1}$ [45, 57]). This decreases, respectively, the coefficients at the efficient members in Eqs. (3b) and (3c) and increases the overall coefficient at the factors of the losses k_{wP870} and k_{wCTC} . The final result is:

$$k_{D,A} = (t_{D,A})^{-1} \approx (1 \pm 0.2) 30.2 \cdot 10^9 \text{ s}^{-1} (4c)$$

and the reverse constant on consideration of the exciton effect is:

$$\delta 16 k_m = 0.35 \cdot 16 (t_m)^{-1} \approx 170 \cdot 10^9 \text{ s}^{-1}. (4d)$$

The system of Eqs. (3, a-f) was solved using the RKGS computer program developed by A. V. Vinogradov [58]. Values of constants of reverse transfers (k_{CTC-} , k_{e1-} , k_{e2-}) were taken from the works of Shuvalov [57] and of Holzwarth and Muller [59]. Kinetics were obtained for the SEE in B875 and all states in the RC from the moment of excitation up to 1.5 ns, and the SEE quantity that finally occurred in the efficient channel and the channel of losses. They correspond to the yields as follows:

Effective quantum yield QY(ph) ≈ 0.887 ;

Quantum yields of losses:

in the B875* antenna complex: **QY(w,A) ≈ 0.086 ;**

in the special P870* pair: **QY(w,P) ≈ 0.027 ;**

total losses: **QY(Sw) ≈ 0.113 .**

To assess the efficiency of the energy migration organization in purple bacteria, values of quantum yields were calculated under conditions of conventional decrease in the distance between B875 molecules and P870 pairs of RCs. The decrease was modeled by the growth of the SEE transfer constants between B875 and P870 on retention of their ratio. Even a 10-fold increase in the rates of k_m and $\delta 16 k_m$ resulted in a rather weak gain in the effective quantum yield value: $0.887 \rightarrow 0.892$. Thus, the SEE migration is organized close to optimal for the antenna of 32 BChl molecules per RC and interchromophore distances between BChl in the LH1 complexes *in vivo*.

DISCUSSION

As shown above, the main quantum losses occur in the BChl antenna fraction B875. Analysis of dependence of these parameters on variations of values of the acting constants revealed that the quantum losses are mainly determined by the constant of the losses in B875, k_{wA} , and the ratio of the forward and reverse migration constants between B875 and P870.

Excluding from consideration of ultrapure crystals, among numerous data of physics and photochemistry on intramolecular migration of energy no precedents are known of efficient deliveries of SEE from large masses of molecules onto a single center. The efficient deliveries are realized only in natural photosynthesis due to the unique feature of membrane proteins, which produce unique *in vivo* cellular structures. Within such structures chlorophyll molecules at a concentration lower than molar are arranged in a manner virtually preventing immediate contacts of their chromophores (!). But the mutual positions of neighboring chromophores promote the domination of long-wavelength bands optimal for the SEE migration in the spectra of absorption bands split by the electromagnetic interaction. These unique properties of transmembrane protein chains provide high rates of migration under conditions of virtual absence of quenching dimers in the mass of antenna chlorophylls. Note that in liquid and solid solutions of chlorophylls the production of quenching dimers begins already at the concentration of about 1–3 mM. In purple bacteria the distances between the chromophore centers in the P870 and B875 complexes are 7.5 and 9 Å [26], which corresponds to local concentrations of this pigment of about 0.5 M. At such distances the migration acts occur with rate more than 10^{12} s^{-1} (!) and the chromophore interactions in the neighboring BChl is sufficiently strong for generation of excitons.

The structure of antenna BChl in photosynthetic membranes organized by the intramembrane parts of proteins provides effective delivery of SEE from the antenna BChl to the reaction centers. This structure is a basis of the effective energetics in photosynthesizing organisms and the most significant achievement in the evolution of living Nature on Earth.

And a complex of fundamental scientific and technological problems emerges in the field of applied solar energetics: to invent something similar to a multilayer photoelectric film from cheap stable polymers and dyes with infrared absorption boundary based on achievements of modern nanotechnology.

In the monograph *The Nanotechnology Revolution* by Drexler and Peterson, the authors compare the level of our modern technologies with the “industrial basis” of natural photosynthesis. They write [60]:

“Industry-as-we-know-it takes things from nature—ore from mountains, trees from forests—and coerces them into forms that someone considers useful. Trees become lumber, then houses. Mountains become rubble, then molten iron, then steel, then cars. Sand becomes a purified gas, then silicon, then chips. And so it goes. Each process is crude, based on cutting, stirring, baking, spraying, etching, grinding, and the like.

Trees, though, are not crude: To make wood and leaves, they neither cut, grind, stir, bake, spray, etch, nor grind. Instead, they gather solar energy using molecular

electronic devices, the photosynthetic reaction centers of chloroplasts. They use that energy to drive molecular machines—active devices with moving parts of precise, molecular structure—which process carbon dioxide and water into oxygen and molecular building blocks. They use other molecular machines to join these molecular building blocks to form roots, trunks, branches, twigs, solar collectors, and more molecular machinery.

Every tree makes leaves, and each leaf is more sophisticated than a spacecraft, more finely patterned than the latest chip from Silicon Valley. They do all this without noise, heat, toxic fumes, or human labor, and they consume pollutants as they go. Viewed this way, trees are high technology. Chips and rockets aren't”.

It is difficult to argue against this conclusion.

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